

Diffraction-grating VCN interferometry and Its Applications

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Workshop on

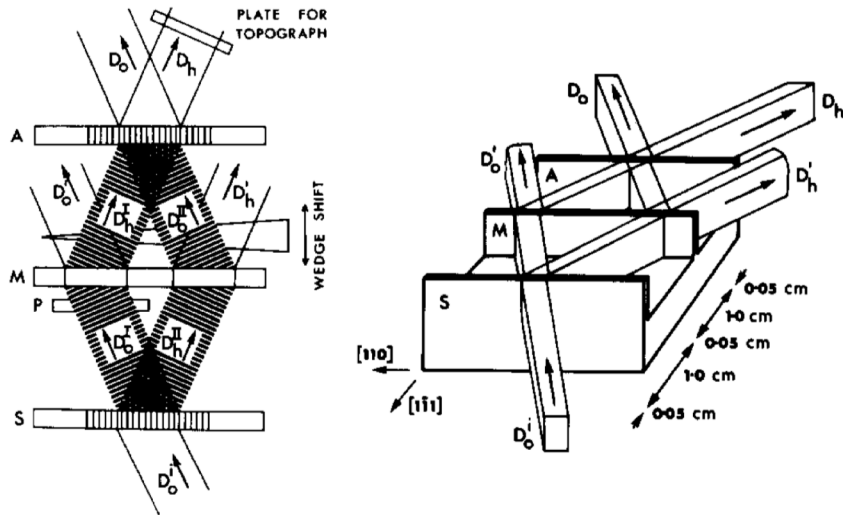
“UCN/VCN sources at the Institute of Nuclear Physics (Kazakhstan) and their applications“

Contents:

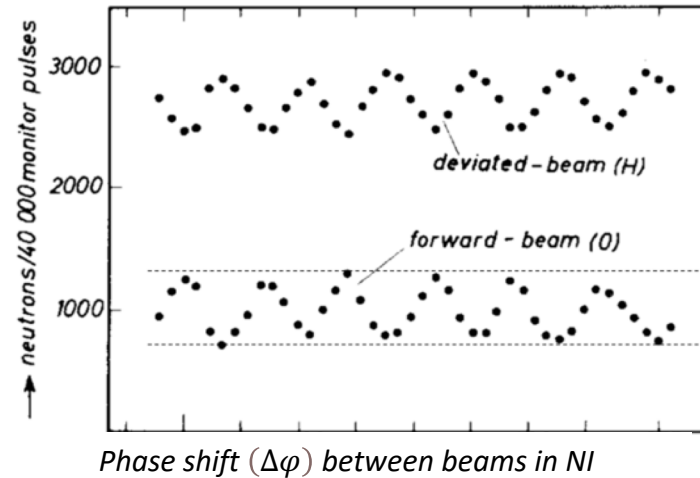
1. VCN diffraction grating interferometer.
2. Applications to neutron charge quest – lowering the experimental limit by orders of magnitude.
3. High-resolution Neutron Speed-Echo spectroscopy of VCN.

Perfect crystal neutron interferometer

Perfect crystal neutron LLL interferometer



Introduced phase shift ($\Delta\varphi$) between beams in the NI

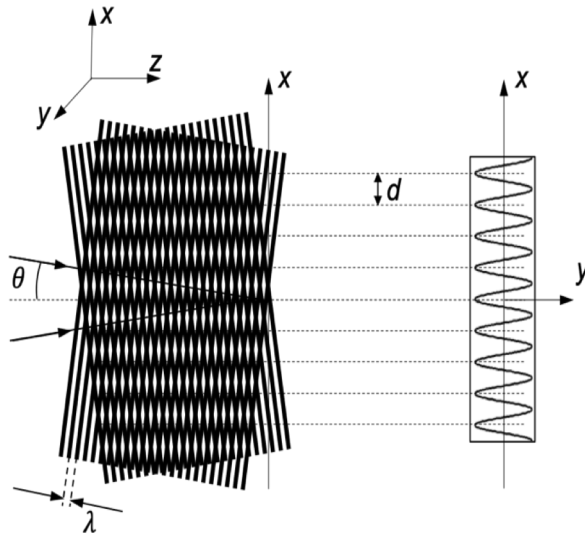


$$I_H(\Delta\varphi) = B - A \cos(\Delta\varphi)$$

Swapping intensity
between O- and H-beams

$$I_O(\Delta\varphi) = V(1 + \cos(\Delta\varphi))$$

H. Rauch et al., 1974



Alternatively:

- crossing coherent waves produce interference pattern of period d
- these interference fringes are superimposed with crystal lattice (d) => Moiré effect (fringes)
- $\Delta\varphi$ results in the lateral shift of interference fringes w.r.t. crystal lattice
- this leads to oscillations of Moiré fringes, i.e. oscillations of recorded intensity

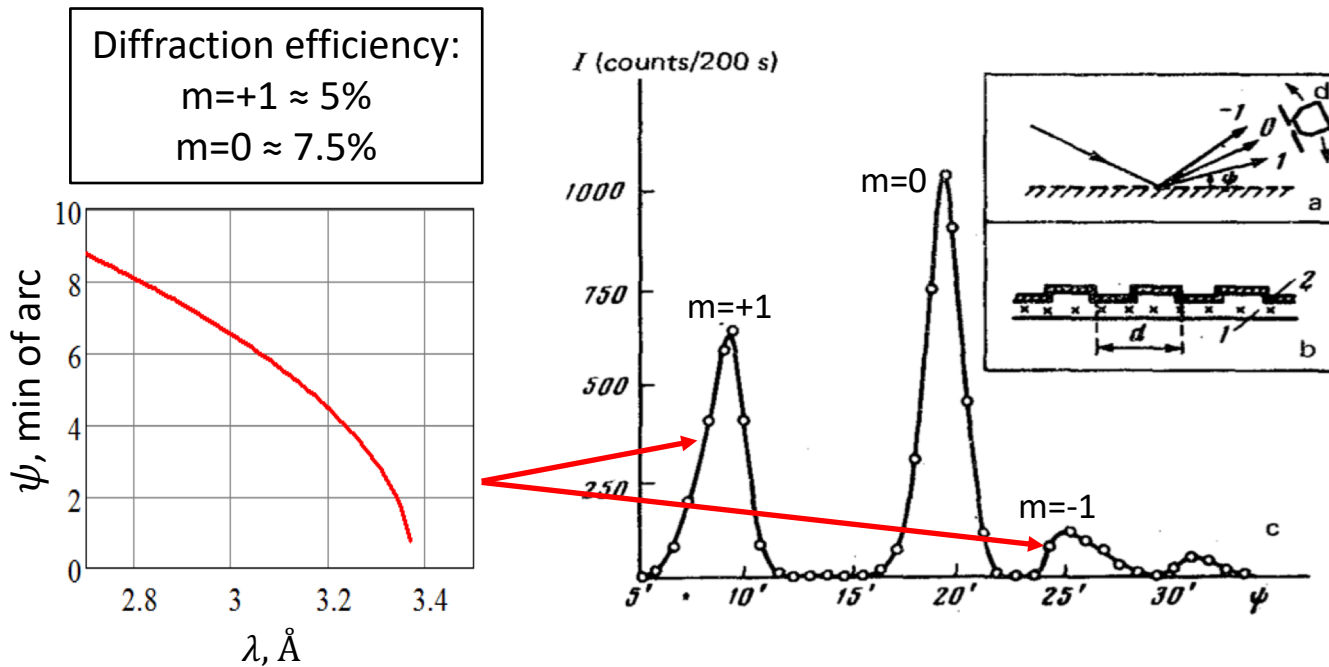
(Very) cold neutrons: coherent beam splitting

For cold neutrons one should employ other than Laue diffraction coherent splitting of neutron waves: diffraction on periodical structures (gratings) or reflection from semi-transparent coatings.

Effective neutron diffraction gratings: modulated surface relief

A.Ioffe et al, JETP letters 33, 374 (1981)

$d = 21 \mu\text{m}$ $\lambda = 2.7 \text{ \AA}$, $\Delta\lambda/\lambda = 32\%$ (Ni mirror)

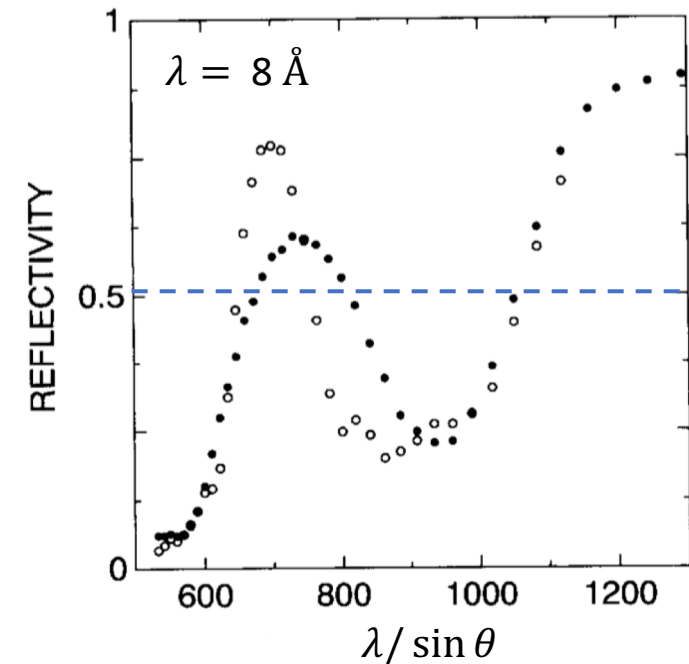


Note asymmetry: spectrum of incident beam => spectroscopy

Coherent beam splitter

T. Ebisawa et al, NIM A 344, 597 (1994)

V-Ti multilayer mirror with spacing $d=360 \text{ \AA}$



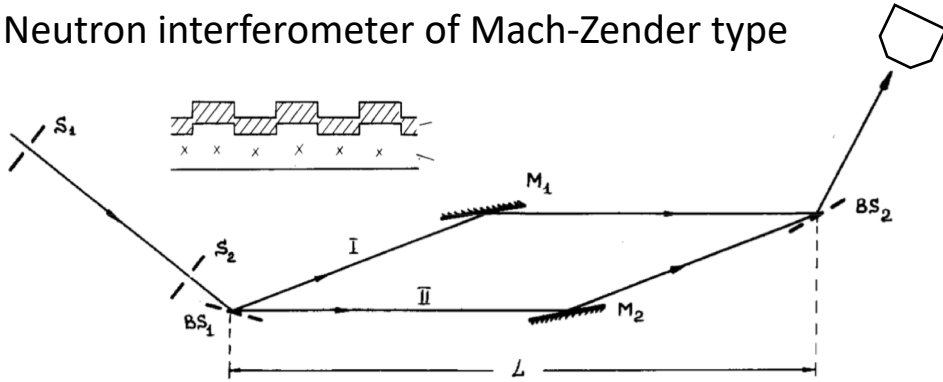
Cold neutron interferometers

Volume 111, number 7 PHYSICS LETTERS 30 September 1985

TEST OF A DIFFRACTION GRATING NEUTRON INTERFEROMETER

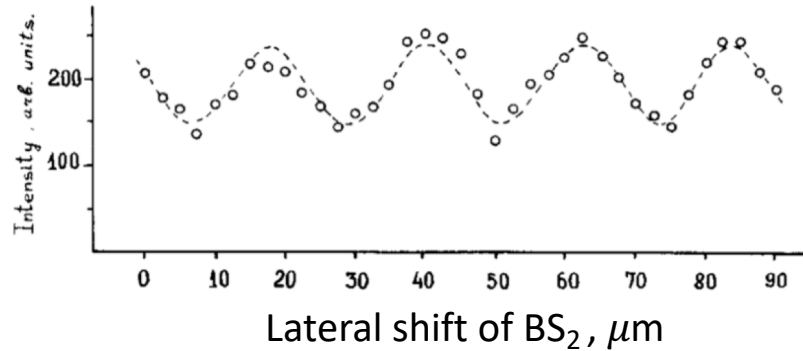
A.I. IOFFE, V.S. ZABIYAKIN and G.M. DRABKIN

Neutron interferometer of Mach-Zender type



Moire pattern with a period d :

$$\lambda = 3.15 \text{ \AA} \quad d = 21 \text{ \mu m}$$



PHYSICAL REVIEW A

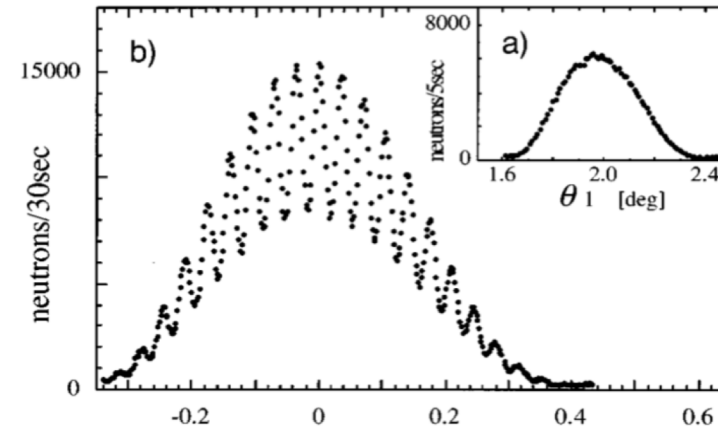
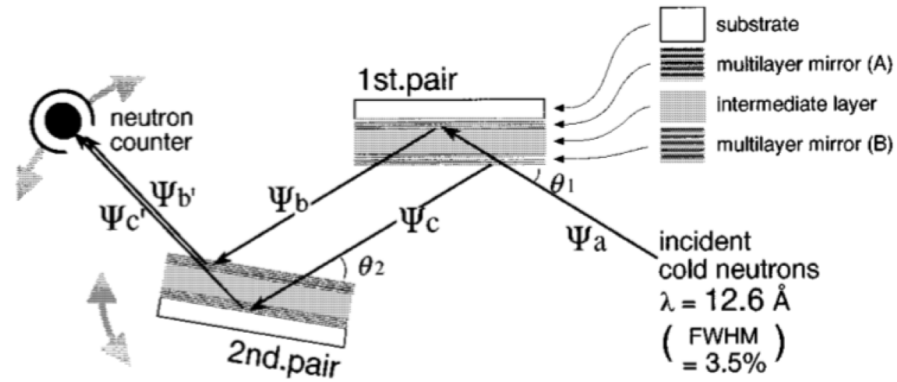
VOLUME 54, NUMBER 1

JULY 1996

Interferometer for cold neutrons using multilayer mirrors

Haruhiko Funahashi,^{1,*} Toru Ebisawa,¹ Tomohito Haseyama,² Masahiro Hino,³ Akira Masaike,² Yoshié Otake,⁴

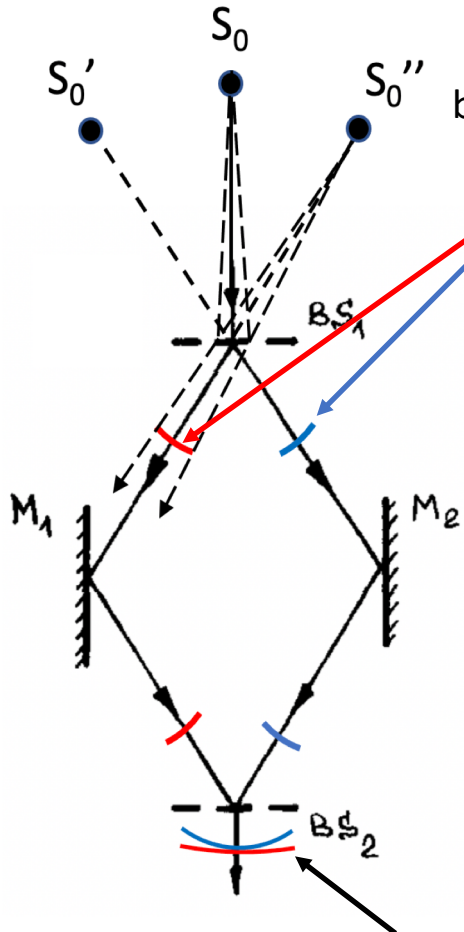
Idea: T. Ebisawa et al, NIM A 344, 597 (1994)



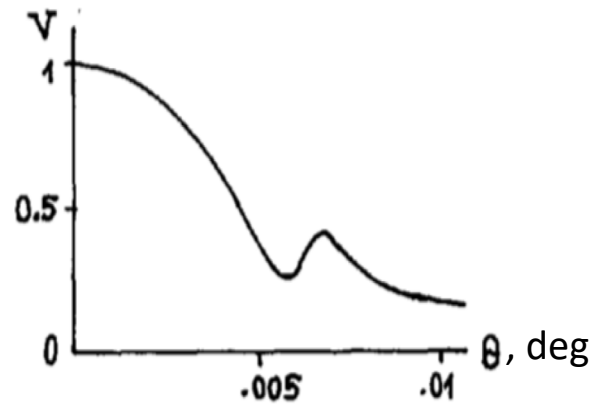
Angle between first and second pairs

See also presentation
of H. Shimizu
at 1st UCN/VCN workshop

3- grating interferometers



Spherical incident wavefronts diffracted by periodic structures are principally aberrated, and non-identical for $m=1$ and $m=-1$.



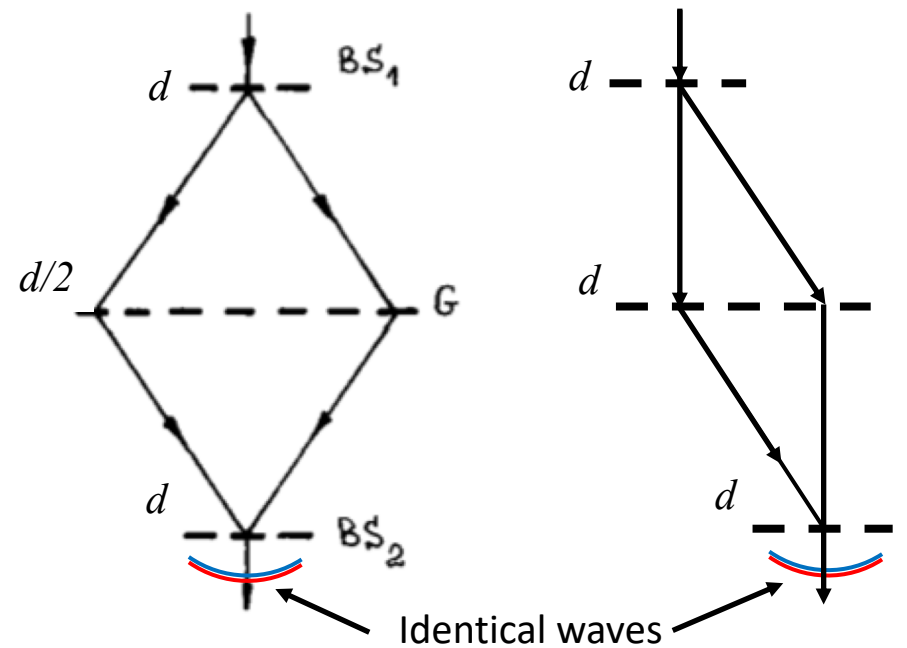
=> Strong requirements to incident beam divergence
=> Bad for neutrons in general; unfeasible for VCN

Interference of two non-identical waves:
=> non-constant period of the interference pattern
=> amplitude modulation over the beam cross-section
=> low visibility V

Unavoidable different aberrations in interfering beams:
=> add complimenting aberrations for equalization.

Deflection --> Diffraction: gratings instead of mirrors

3-grating interferometers

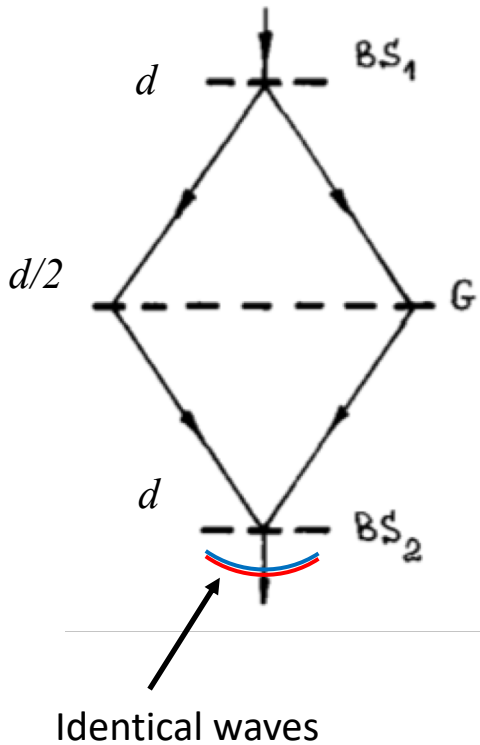


Aberration analysis shows that now interfering wavefronts are distorted identically and $V=1$:
=> no requirements to incident beam divergence

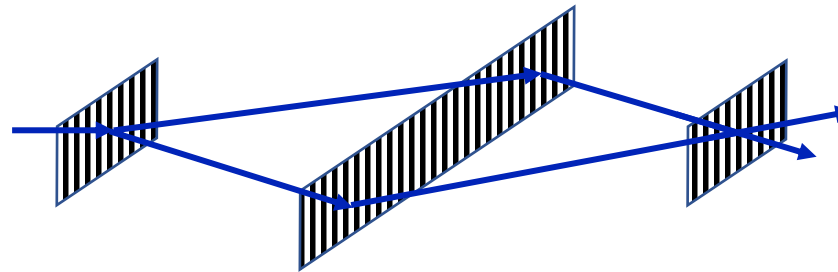
Diffraction grating interferometers

This is not the Talbot interferometer: Talbot effect is a near-field diffraction effect, where the self-imaging of periodic objects (gratings) **requires spatially coherent illumination**.

Here: the imaging of a grating by a second grating **regardless of the coherence of the source**.



- First shown by first-order diffraction theory (i.e. without accounting for aberrations): (*B.Chang, R.Alferness, E.Leith (Appl. Opt. 14 (1975) 1569)* .
- Aberration analysis (higher-orders diffraction theory): full compensation of aberrations => interfering waves are identical (*A.Ioffe, NIM A268 (1988) 169*).



Such interferometer works regardless of the source coherence, i.e. for non-monochromatic and non-collimated neutron beam!

Transition to neutrons:

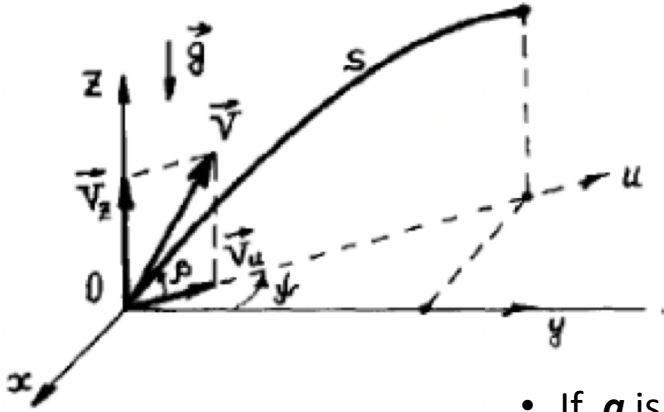
refraction index of vacuum in gravitational field $\neq 1$.

(*I.M Frank, A.I Frank, JETP Lett. 28 (1978) 515*)

$$n = \sqrt{1 - 2gz \left(\frac{m_n}{h}\right)^2 \lambda^2} \Rightarrow$$

As neutrons propagate on parabolic trajectories, vacuum has non-linear refraction index. This is not trivial, will be discussed later.

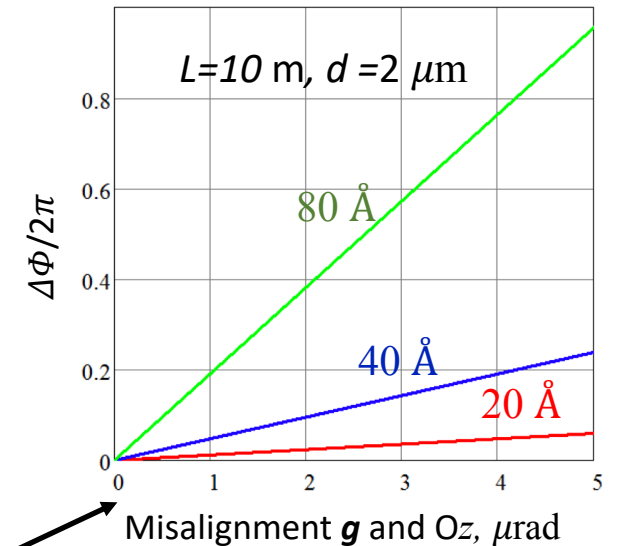
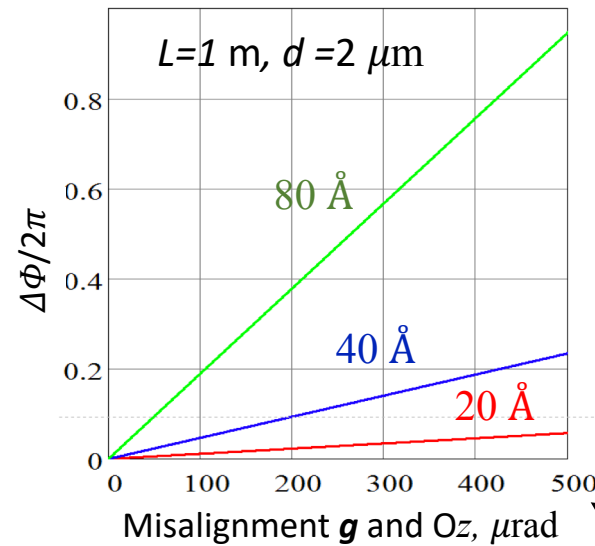
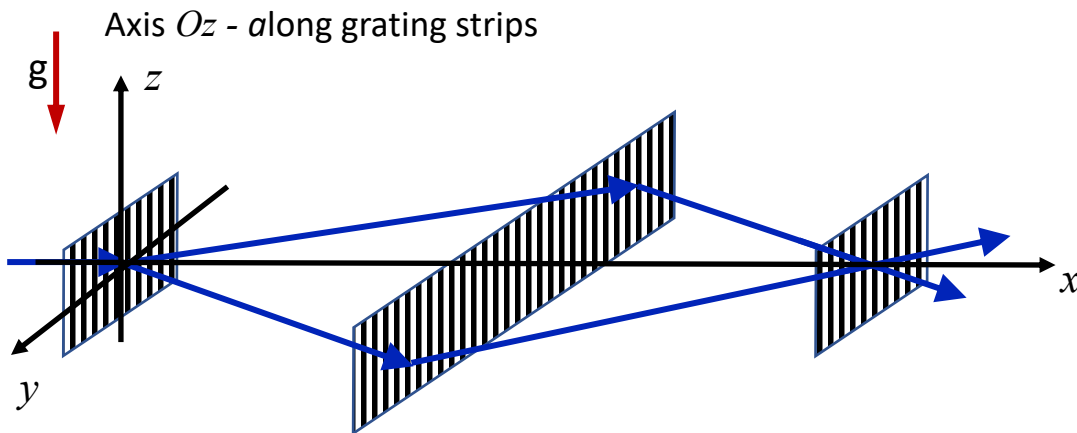
VCN grating interferometers in gravitational field



$$\Phi(u) = \frac{m_n}{\hbar} \left[uV_u + \frac{V_z^3}{3g} - \frac{V_u^3}{3g} \left(\tan \beta - \frac{gu}{V_u^2} \right)^3 \right],$$

Calculating the velocity components immediately after diffraction, it is possible to calculate the phase shift of neutron wave during its following propagation.

- If g is strictly parallel to Oz (grating strips), gravitational potentials for both sub-beams are equal (symmetry).
- Violation of this symmetry leads to neutron trajectories rising to different heights => phase difference.



Note different scales!

3-grating interferometer

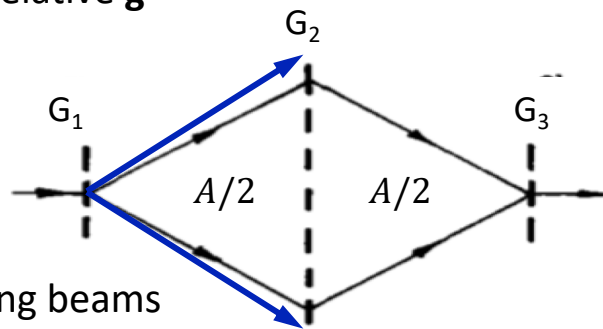
(+) for VCN: aberration-free, V=100% for full incoherent illumination

(-) for VCN: requires μrad alignment relative \mathbf{g}

(-) parasitic Sagnac effect

$$\varphi_S = \frac{2m_n}{\hbar} (\boldsymbol{\omega} \cdot \mathbf{A}) = \frac{2m_n}{\hbar} \omega_0 A \sin \theta_1,$$

$$A = \frac{\lambda L}{d^2} \text{ - area enclosed by interfering beams}$$



Problem is not φ_S , rather $\Delta\varphi_S$ because of different A

Scatter in λ : $\Delta\lambda \rightarrow \Delta A \rightarrow \Delta\varphi_S$ $\Delta\varphi_S = \varphi_S \frac{\Delta\lambda}{\lambda} = 0.1A$

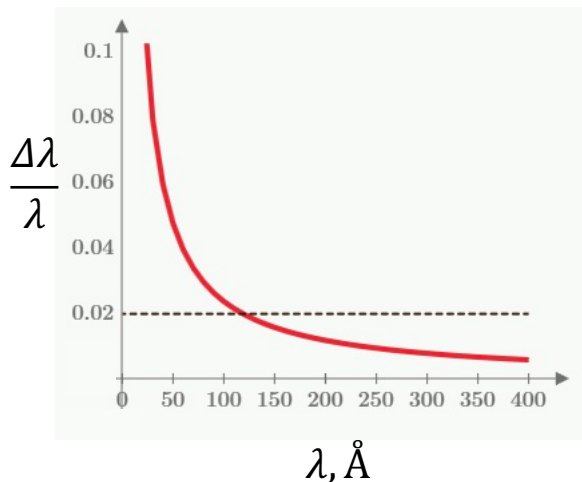
$\omega_0 = 7.29 \cdot 10^{-5} \text{ s}^{-1}$ angular velocity (Earth's rotation)

$\theta_1 \approx 43^\circ$ - latitude angle (Almaty)

Interference fringes are washed out for $\Delta\varphi_S > \pi/2$:

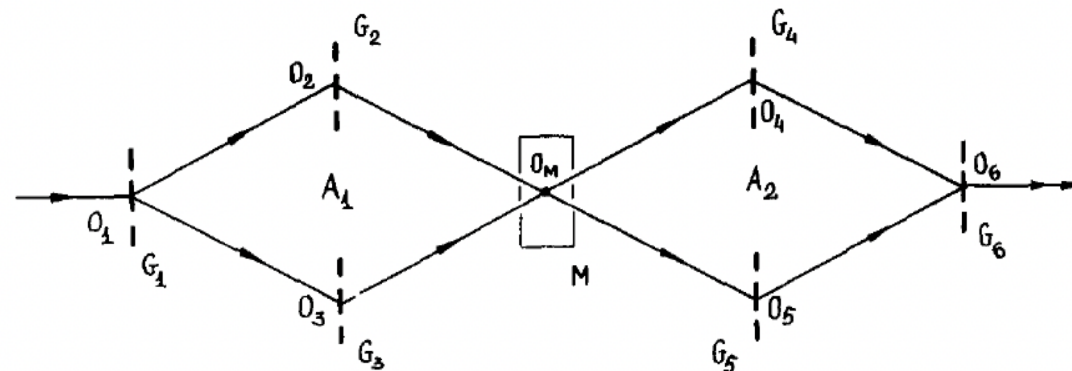
=> limitation on $\frac{\Delta\lambda}{\lambda}$

Not feasible with VCNs.



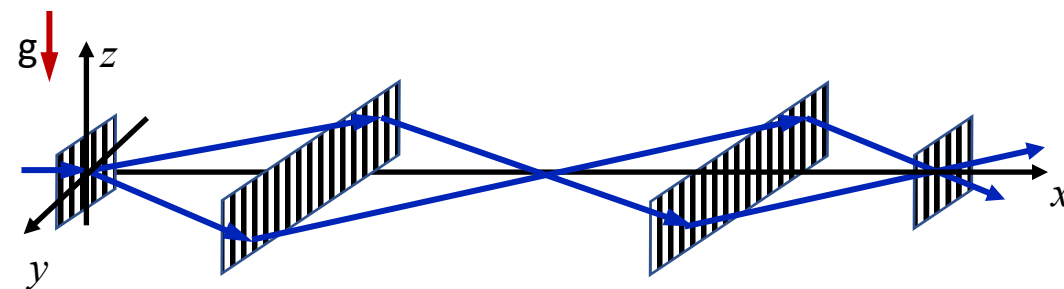
4-grating interferometer

A.Ioffe, NIM A268 (1988) 169.



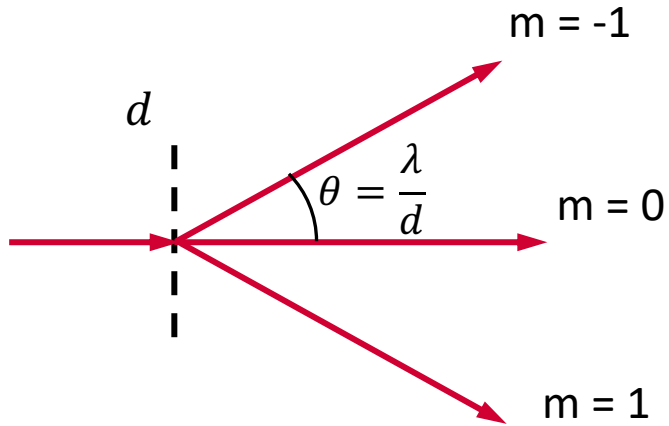
Symmetric scheme:

- $A_1 = -A_2$ - complete compensation of Sagnac effect
- complete compensation of gravitational phase difference



Not for free: one more grating - additional intensity losses

Diffraction gratings for VCNs

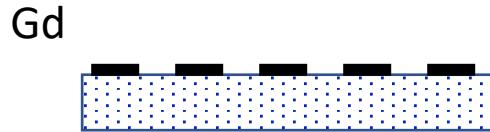


Requirements: small period and high diffraction efficiency

- Photolithographic gratings (stamping in photoresist)
- Holographic photolithographic gratings (interference lithography)
- Holographic nanodiamond-polymer composite gratings (E.Hadden et al, Appl. Phys. Lett. 124, 071901 (2024))

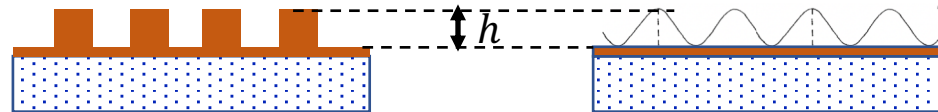
Phase gratings

Amplitude gratings



Phase shift: $\varphi = \frac{2\pi\lambda}{\rho} h$

ρ - scattering length density
 $h = 1.74 \mu\text{m}$ for $\lambda = 20 \text{ \AA}$.



Diffraction efficiency:

$$\eta_m = \frac{1}{\pi^2 m^2}$$

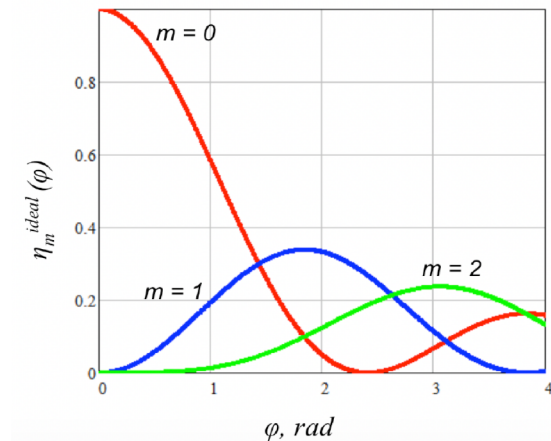
$$\eta_1^{\text{max}} = 10.1\%$$

$$\eta_m = \frac{4}{\pi^2 m^2} \sin^2(\varphi)$$

$$\eta_1^{\text{max}} = 40.4\%$$

$$\eta_m = [J_m(\varphi)]^2$$

$$\eta_1^{\text{max}} = 33.8\%$$



First realization of VCN diffraction grating interferometer

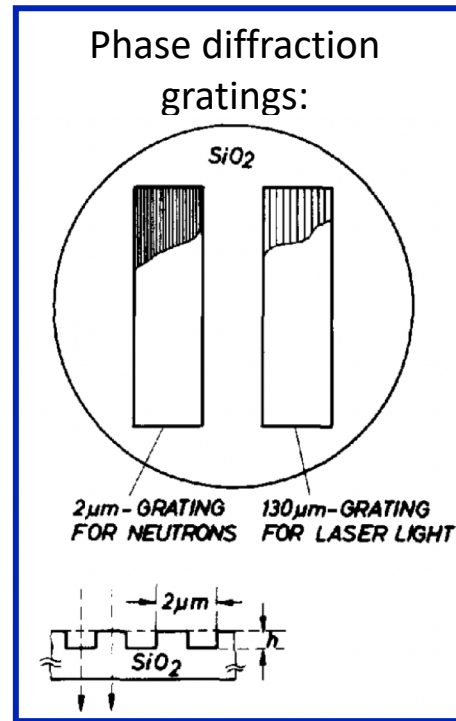
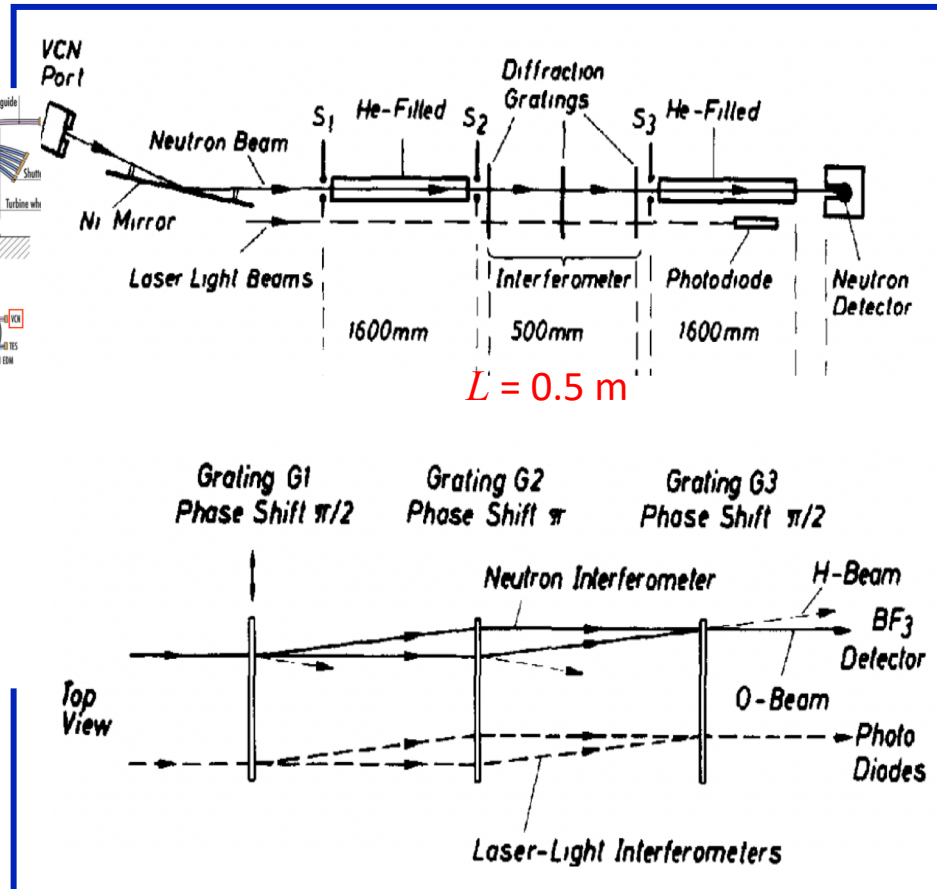
Volume 140, number 7,8

PHYSICS LETTERS A

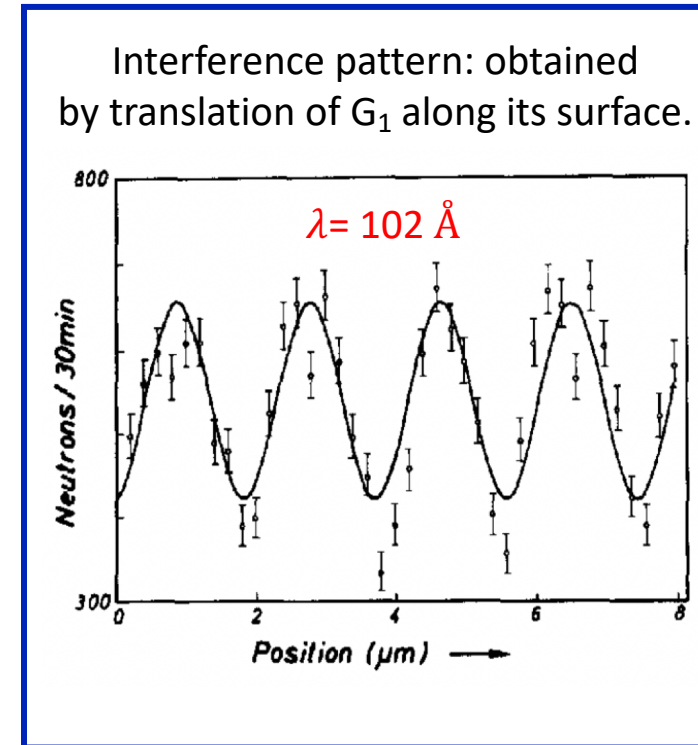
9 October 1989

A PHASE-GRATING INTERFEROMETER FOR VERY COLD NEUTRONS

M.Gruber, K.Eder, A.Zeilinger, R.Gähler, W.Mampe



$d = 2 \mu\text{m}$



Potential application of VCN interferometry: searching for the net electric charge of neutrons

- I will not discuss “why” to measure q_n , just “how” to measure it.
- Earlier attempts and current experimental limit on q_n .
- VCN interferometry
- Utilizing VCN to improve the experimental limit on q_n .

Previous experiments and limits on neutron charge q_n

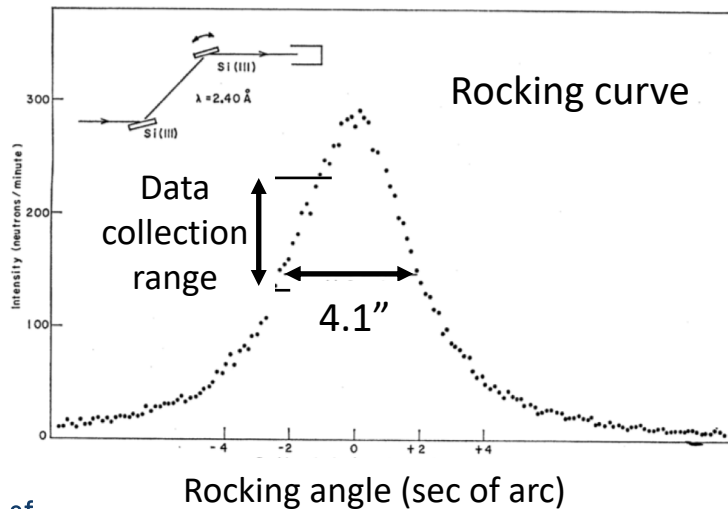
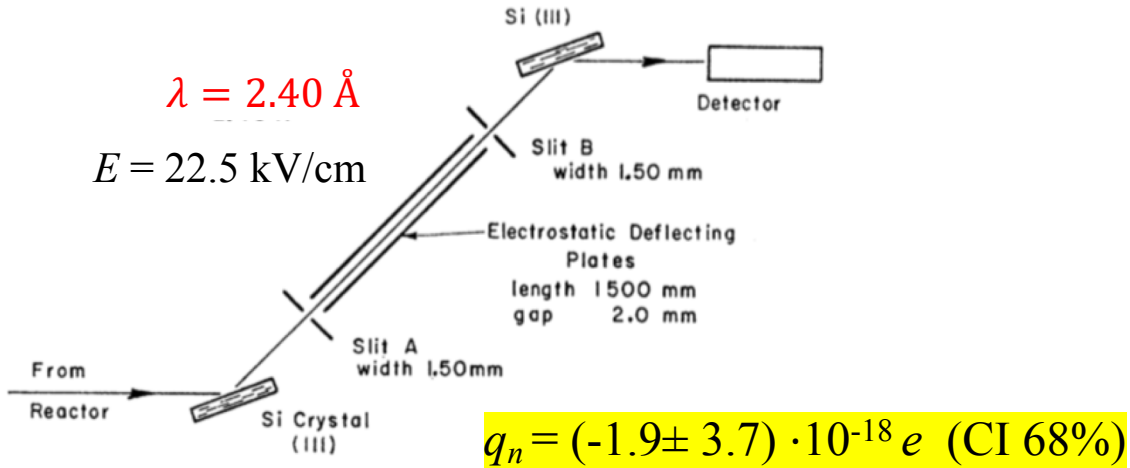
PHYSICAL REVIEW

VOLUME 153, NUMBER 5

25 JANUARY 1967

Experimental Limit for the Neutron Charge*

C. G. SHULL, K. W. BILLMAN, AND F. A. WEDGWOOD†



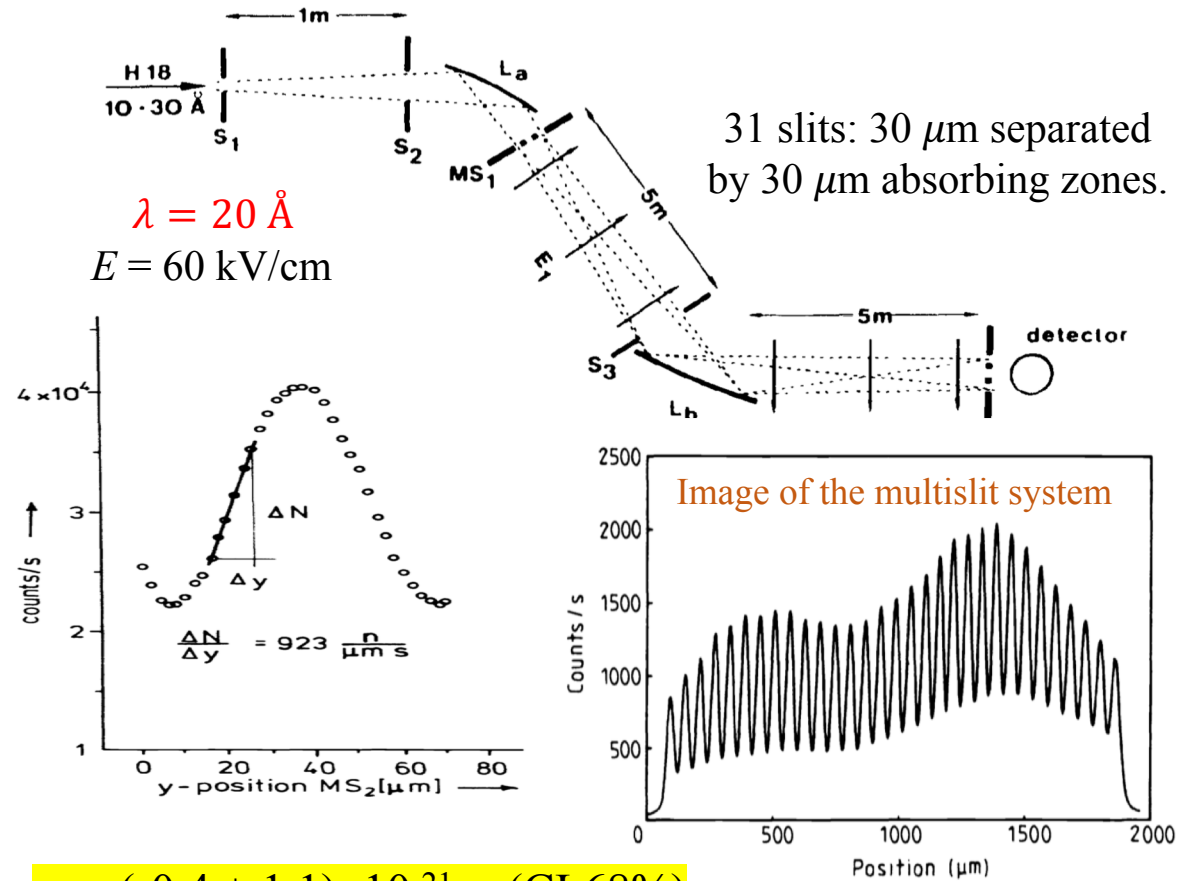
PHYSICAL REVIEW D

VOLUME 25, NUMBER 11

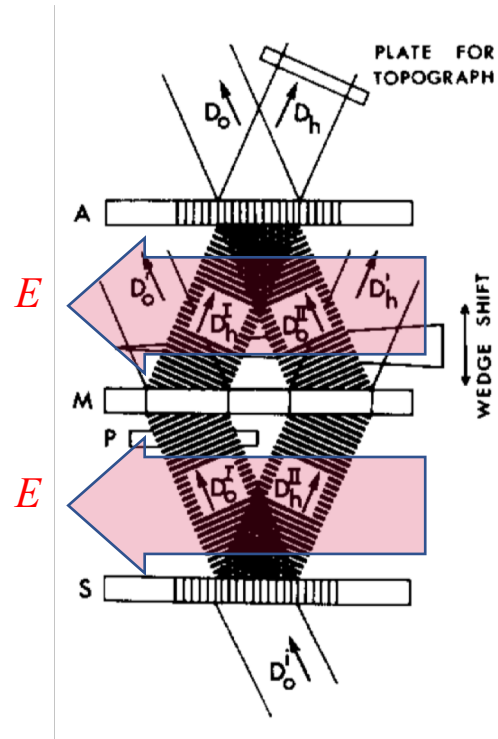
1 JUNE 1982

Experimental limit for the charge of the free neutron

R. Gähler, J. Kalus, W. Mampe



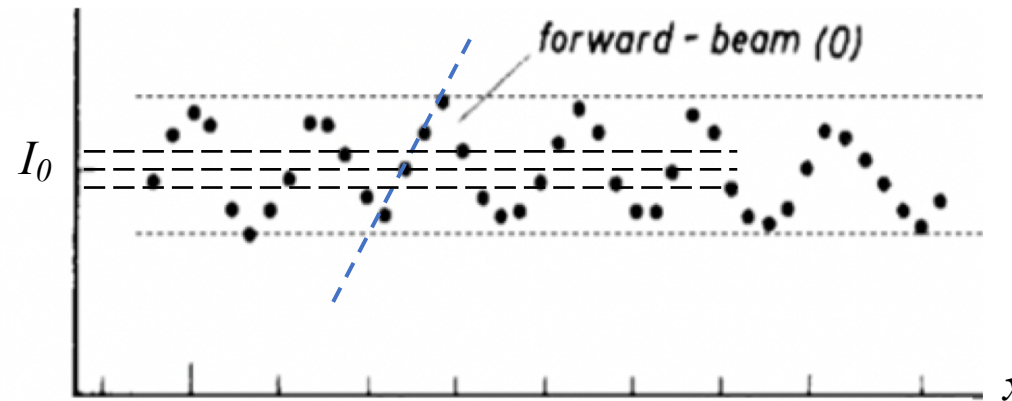
Gedanken experiment with crystal interferometer: neutron charge q_n



Electric field across neutron beams => shift of interference pattern:

$$\Delta x = \frac{1}{2} qE \left(\frac{L}{v}\right)^2 = \frac{1}{2} qE \left(\frac{L}{h} m\lambda\right)^2 \sim qEL^2\lambda^2$$

$$I(\Delta x) = I_0 V \left(1 + \cos \frac{2\pi}{d} \Delta x\right) \quad \rightarrow \quad \text{Shift by } d \text{ (lattice spacing) corresponds to full intensity oscillation}$$

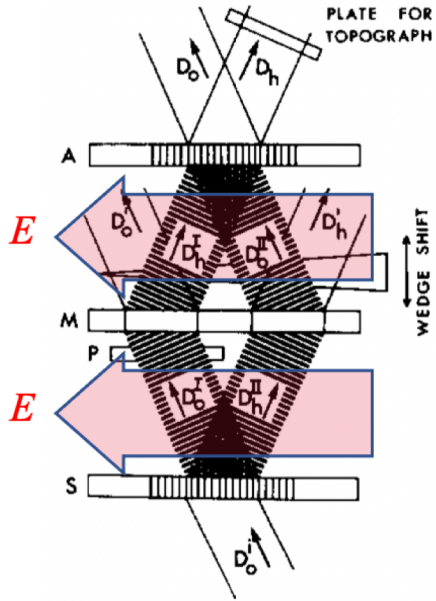


$$q = \frac{\sqrt{2} d}{\pi \sqrt{I_0}} \frac{1.5}{V E L^2 \lambda^2} \left(\frac{h}{m}\right)^2$$

For $E = 60 \text{ kV/cm}$, $L = 5 \text{ cm}$, $\lambda = 2 \text{ \AA}$, $d = 1.92 \text{ \AA}$ $q_n \geq 3 \cdot 10^{-20} e$ (CI 90%) in 100 days

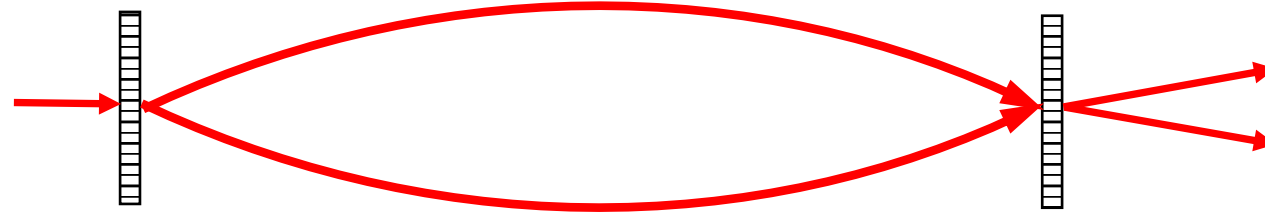
Important: $q \sim \frac{d}{\sqrt{I_0} E L^2 \lambda^2}$ - this is a kind of FOM

Neutron interferometer with larger length and wavelength: VCN



$$q_n \sim \frac{d}{\sqrt{I_0} E L^2 \lambda^2}$$

Now imagine we can modify our interferometer towards larger length and wavelength, however with corresponding increase of d .



Scaling to VCN:

$$\lambda: 2 \text{ \AA} \rightarrow 20 \text{ \AA} \Rightarrow \times 10^2$$

$$L: 5 \text{ cm} \rightarrow 5 \text{ m} \Rightarrow \times 10^4$$

x

$$d: 2 \text{ \AA} \rightarrow 1 \text{ \mu m} \Rightarrow \times (2 \cdot 10^{-4})$$

$$I_0: \sqrt{\lambda^{-5}} \Rightarrow \times (3 \cdot 10^{-3})$$

=>

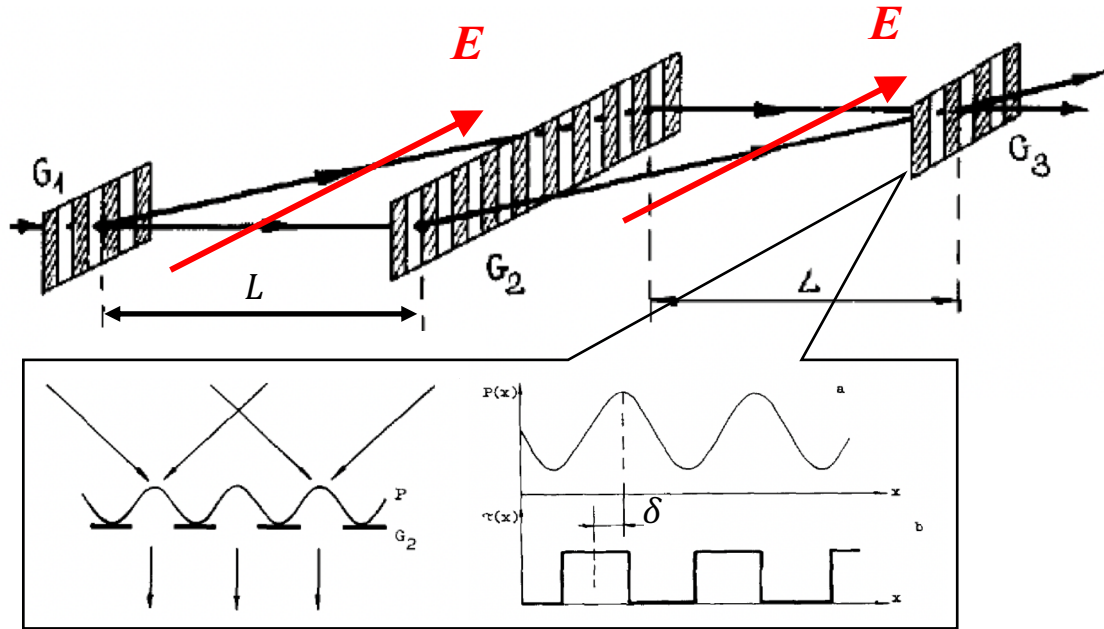
Thermal to cold
neutron source: x15

Total gain about 10 => one can put a harder limit on q_n

However, for cold neutrons one should use other than Laue diffraction coherent splitting of neutron waves.

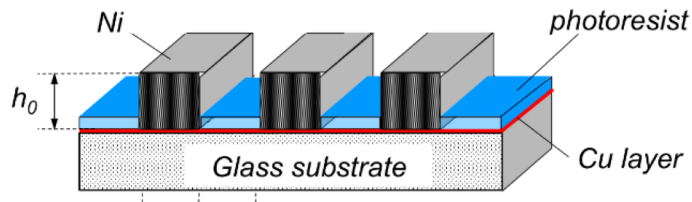
VCN diffraction grating interferometer for search of q_n

Electric field applied across interferometer beams



$$I(\Delta x) = I_0 V \left(1 + \cos \frac{2\pi}{d} \delta \right)$$

Phase diffraction gratings: surface relief



$d = 3.3 \mu\text{m}$
 $h_0 = 1.7 \mu\text{m}$: phase shift π for $\lambda = 20 \text{ \AA}$

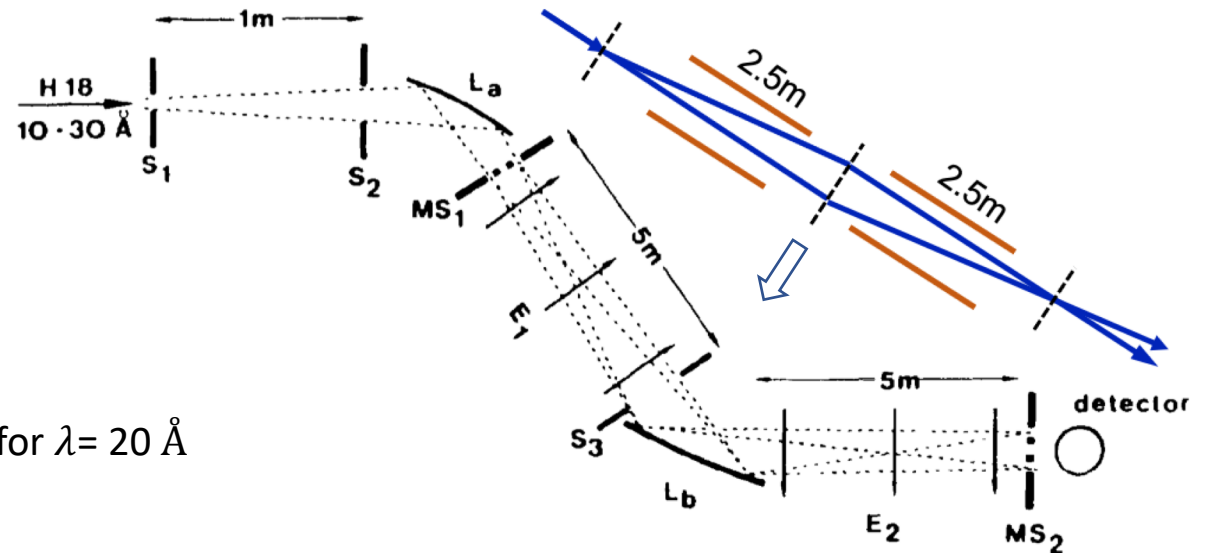
A. Ioffe, NIM A228 (1984) 141; NIM A268 (1988) 169.

$$\delta = \frac{1}{2} q E \left(\frac{L}{h} m \lambda \right)^2 \quad q = \frac{\sqrt{2} d}{\pi \sqrt{I_0}} \frac{1.5}{E L^2 \lambda^2} \left(\frac{h}{m} \right)^2$$

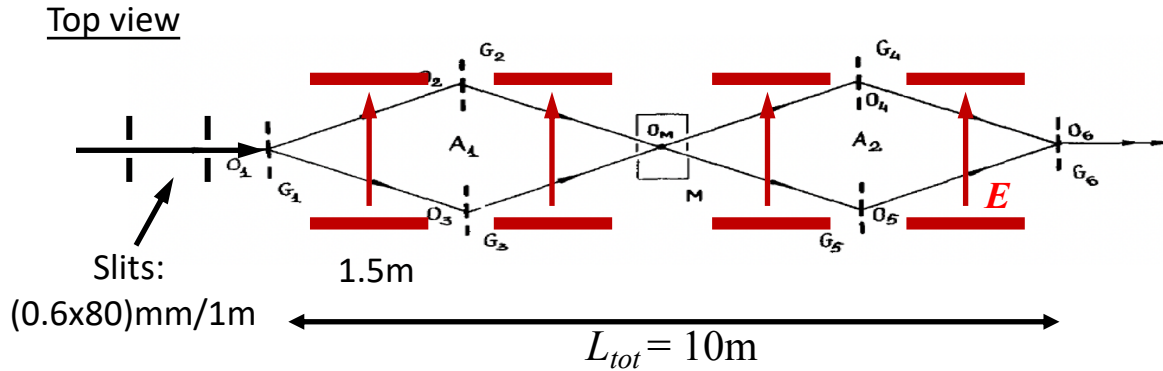
1987- proposal to ILL (accepted, but was not materialized):
 to use the same setup at H18 as for previous q_n experiment:

$I_0 = 200 \text{ n/s}$, $\lambda = (20 \pm 0.15) \text{ \AA}$, $E = 60 \text{ kV/cm}$, $L = 5 \text{ m}$

$q_n \geq 2 \cdot 10^{-22} e$ in 60 days - order of magnitude improvement



VCN diffraction grating interferometer at ESS: search for q_n



$$d = 2 \mu\text{m}$$

$$V=50\%$$

$$E = 60 \text{ kV/cm}$$

$$L_E = 6\text{m}$$

Beam parameters: the same as at the ILL setup

$$\text{Beam cross-section: } S_{\text{beam}} = 0.48 \text{ cm}^2$$

$$\text{Solid angles: } \omega_x = 0.0006$$

$$\omega_y = 0.08 \text{ (L=10m, last slit is } G_4)$$

$$\Delta\lambda = 3956 \cdot T_{\text{rep}} / L_{\text{source-det}} = 14 \text{ \AA}$$

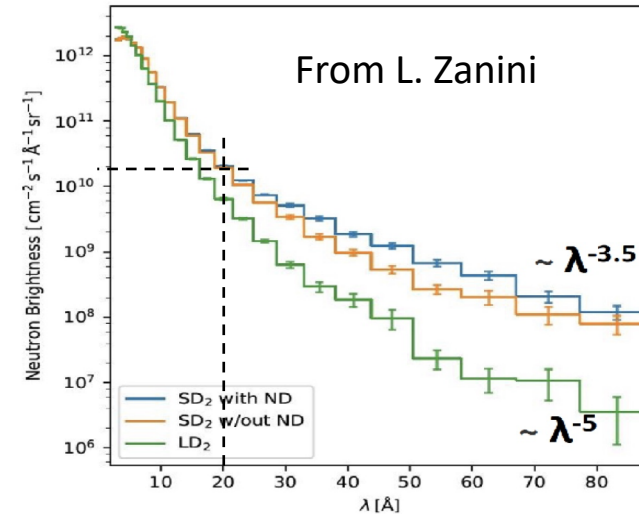
$$\text{Diffraction efficiency: } \eta^4 = 0.008 \text{ (}\eta = 30\%)$$

Transmission of substrates ($\lambda = 20 \text{ \AA}$):

$$\text{Si } 4 \times 0.07 \text{ cm: } T_{\text{Si}} = 0.93$$

$$\text{SiO}_2 \text{ } 4 \times 0.3 \text{ cm: } T_{\text{SiO}_2} = 0.63$$

$$I_{\text{rec}}(\lambda) = B(\lambda) S_{\text{beam}} \omega_x \omega_y \eta^4 \Delta\lambda T_{\text{SiO}_2}$$

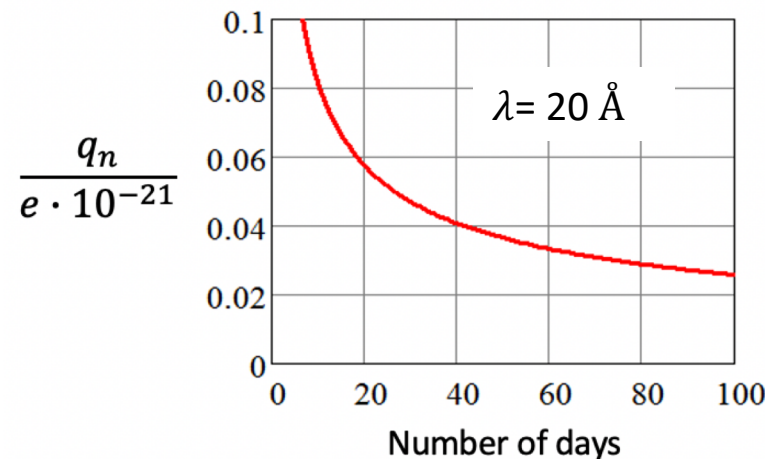


$$B(20 \text{ \AA}) = 2 \times 10^{10}$$

Expected counting rate:

$$I_{\text{rec}}(20 \text{ \AA}) \approx 4.9 \cdot 10^3 \text{ n/s}$$

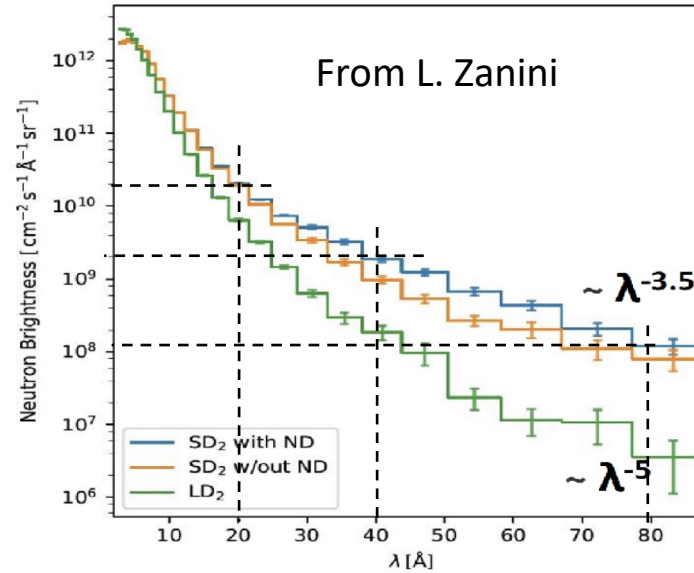
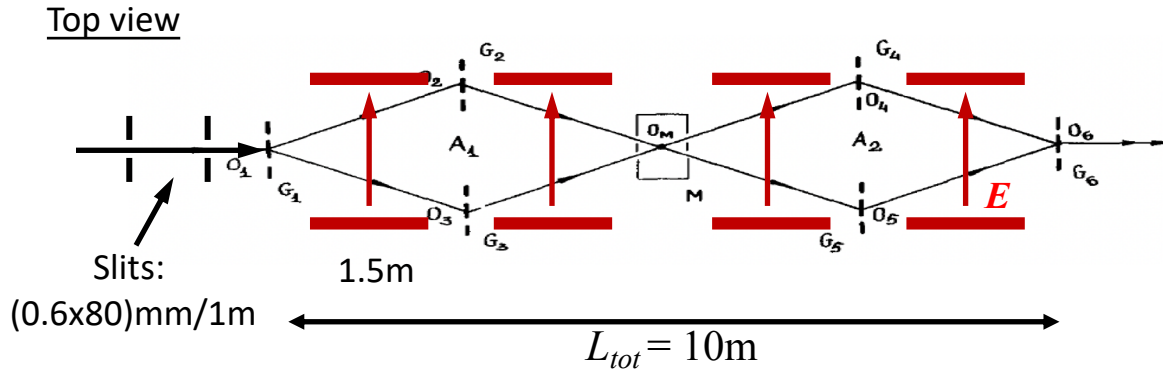
$$q(\lambda) = \frac{\sqrt{2} d}{\pi \sqrt{I_{\text{rec}}(\lambda) 8.64 \cdot 10^4 N_{\text{days}}}} \frac{1.65}{V E L^2 \lambda^2} \left(\frac{h}{m}\right)^2$$



$$q_n \geq 3 \cdot 10^{-23} e \text{ in 80 days (CI 90\%)}$$

2 orders of magnitude better,
than the present day limit

VCN diffraction grating interferometer at ESS: search for q_n



Expected counting rate:

$$I_{rec}(20 \text{ Å}) \approx 4.9 \cdot 10^3 \text{ n/s}$$

$$I_{rec}(40 \text{ Å}) \approx 4.9 \cdot 10^2 \text{ n/s}$$

$$I_{rec}(80 \text{ Å}) \approx 4 \cdot 10^1 \text{ n/s}$$

Transition to higher λ ($> 20 \text{ Å}$): does it make sense?

$$q_n \sim \frac{d}{\sqrt{I_0 E L^2 \lambda^2}}$$

$$\text{LD}_2: q_n \sim \frac{1}{\lambda^{-2.5} \lambda^2} \sim \lambda^{0.5} \quad \text{Getting worse}$$

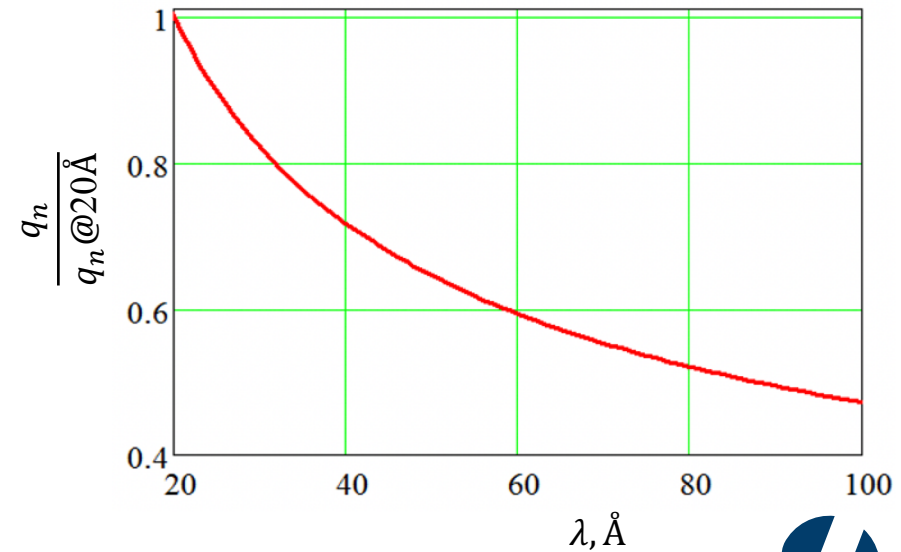
$$\text{SD}_2: q_n \sim \frac{1}{\lambda^{-1.75} \lambda^2} \sim \lambda^{-0.25} \quad \text{Getting better}$$

SD₂ with ND is a game changer:

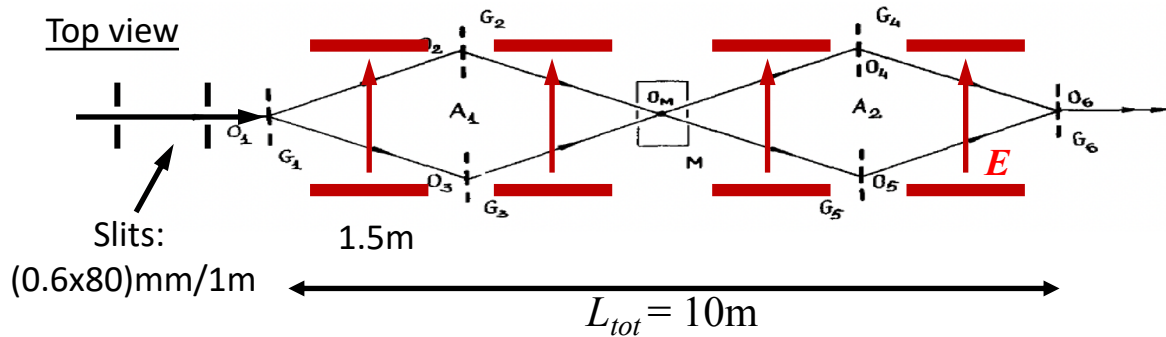
Transition from 20 Å to 80 Å gives factor of 2 improvement:

$$q_n \geq 1.5 \cdot 10^{-23} e \text{ in 80 days (CI 90\%)}$$

Cold source at WWR-K?



Potential for further improvements in search for q_n



$V=50\%$, $E = 60$ kV/cm, $L_E = 6m$

Beam parameters:

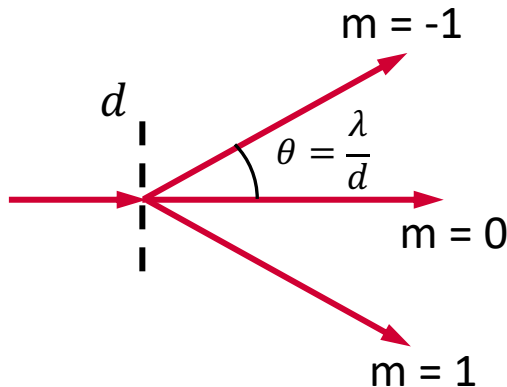
Beam cross-section: $S_{beam} = 0.48$ cm²

Solid angles:

$\omega_x = 0.0006 \ll \lambda/d = 0.005$

$\omega_y = 0.08$ (L=10m, last slit is G₄)

Band: $\Delta\lambda = 14$ Å



$$q_n \sim \frac{d}{\sqrt{I_0} E L^2 \lambda^2}$$

Practically, only “free” parameter is grating period d .

Reducing the period d of diffraction gratings to sub- μ m:

\Rightarrow direct gain as $q_n \sim d$

\Rightarrow increase of diffraction angle $\theta = \frac{\lambda}{d}$,

therefore gain in incident beam intensity $\sim d^2$:

gain in solid angle $\omega_x \sim d$ (still $\ll \lambda/d$)

gain in beam cross-section $\sim d$

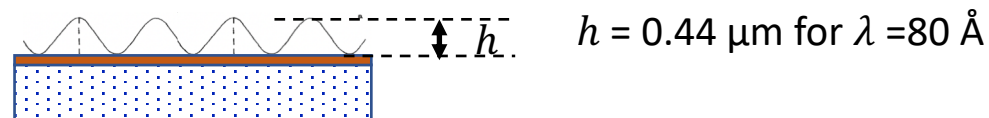
Overall gain in $q_n \sim d^{-2}$

$d: 2 \mu\text{m} \rightarrow 0.5 \mu\text{m}$ results in additional improvement by an order of magnitude: $q_n \geq 10^{-24} e$

\Rightarrow Holographic (interference) gratings

\Rightarrow Holographic ND-composite gratings

(*E.Hadden et al, Appl.Phys.Lett., 2024*)

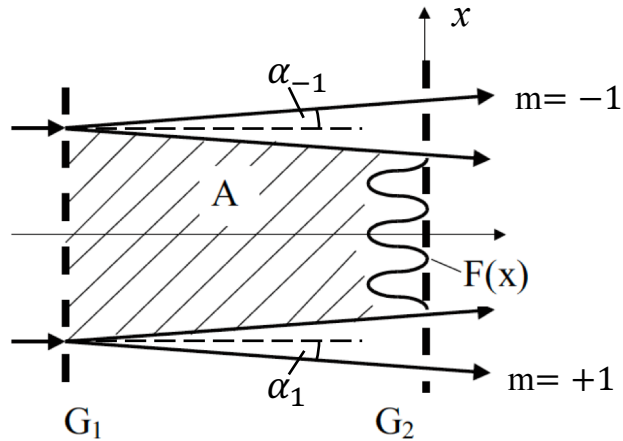


**Another potential application for VCN:
very high resolution Neutron Speed Echo spectroscopy (NSPE)**

Neutron Speed Echo (NSPE): principle

Two stationary diffraction gratings

A. Ioffe, in Neutron Spin Echo Spectroscopy, Lecture Notes in Physics (2002) p.142.



Diffraction on G_1 :

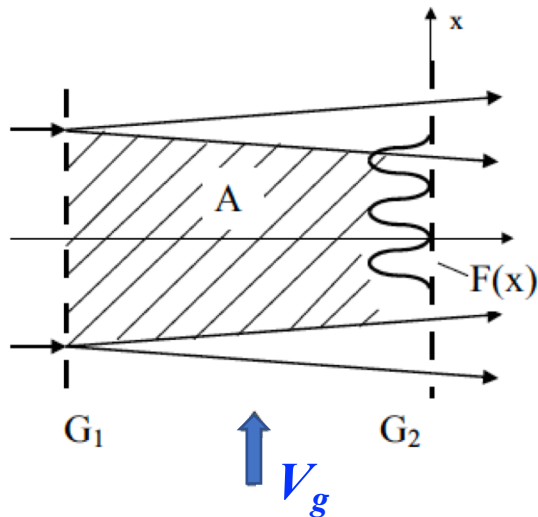
- amplitude division of wavefronts
- diffracted waves are coherent
- in superposition region: sinusoidal amplitude distribution $F(x)$ with period d_F



Overlay G_2 and $F(x)$:
Moiré patterns are the same for all λ

$$d_F = \frac{\lambda \cdot \sin(\alpha_1 + \alpha_{-1})}{2} = d \Rightarrow \text{self-imaging of } G_1 \text{ in } G_2$$

Moving diffraction gratings (velocity V_g)

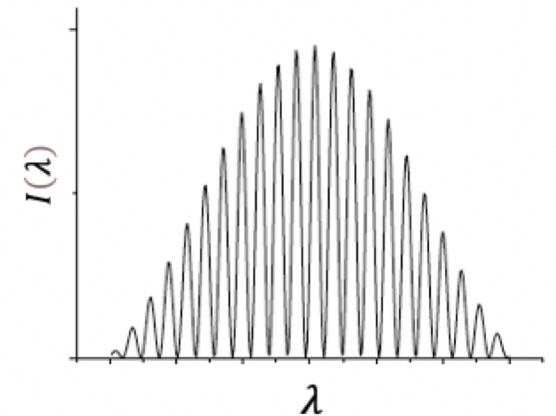


Different neutron wavelength λ :

- \Rightarrow different propagation time between G_1 and G_2
- \Rightarrow Shift of $F(x)$ relative G_2
- \Rightarrow Moiré patterns are different for different λ :

$$I(\lambda) = \frac{I_0(\lambda)}{2} \left\{ 1 + \cos \left[\frac{2\pi}{d} V_g \frac{m_n}{h} L \lambda \right] \right\}$$

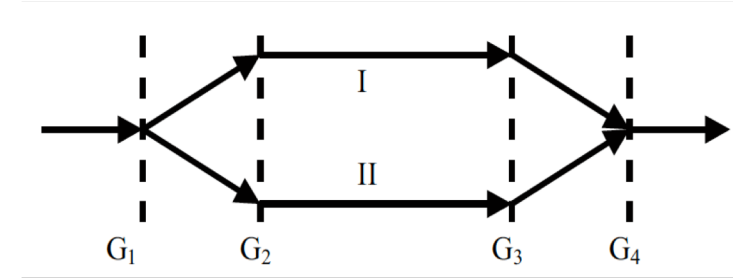
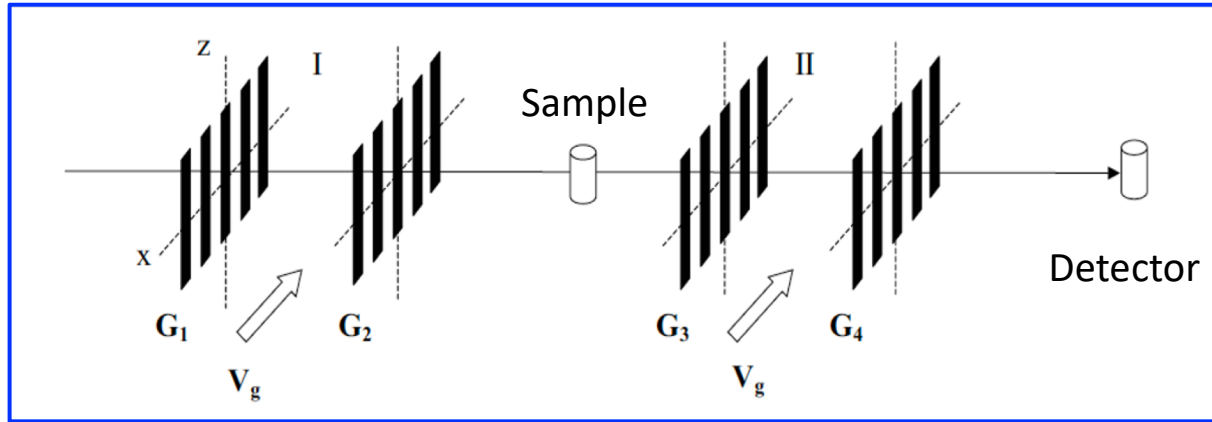
Modulation of outgoing neutron beam spectrum: $f_g = \frac{V_g}{d}$



Neutron Speed Echo spectrometer

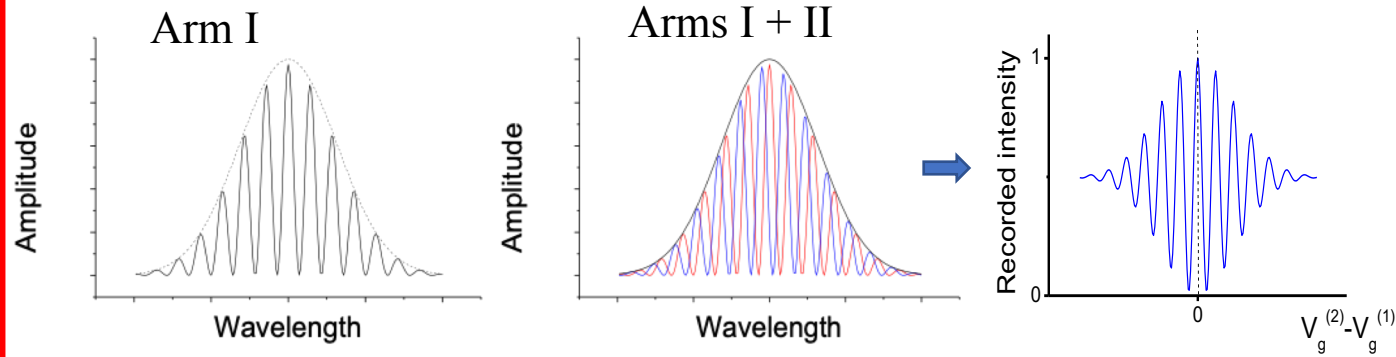
A. Ioffe, Physica B283 (2000) 406.

Two similar arms:



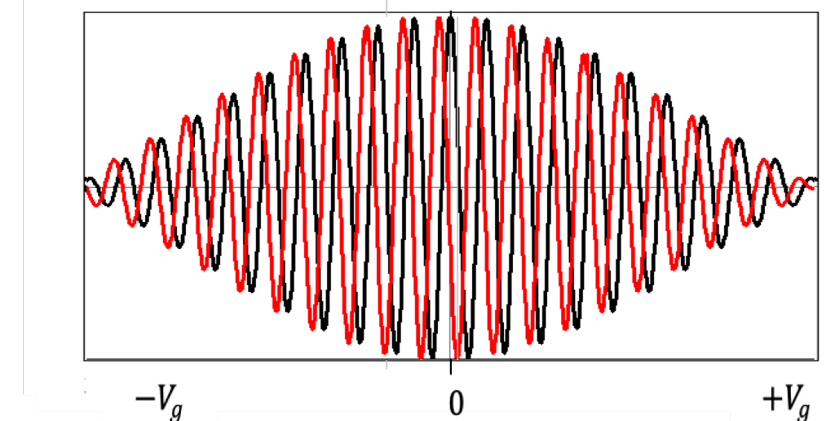
4-grating interferometer: 100% contrast of interference fringes for non-collimated and non-monochromatic neutron beam.

Neutron spectrum modulation

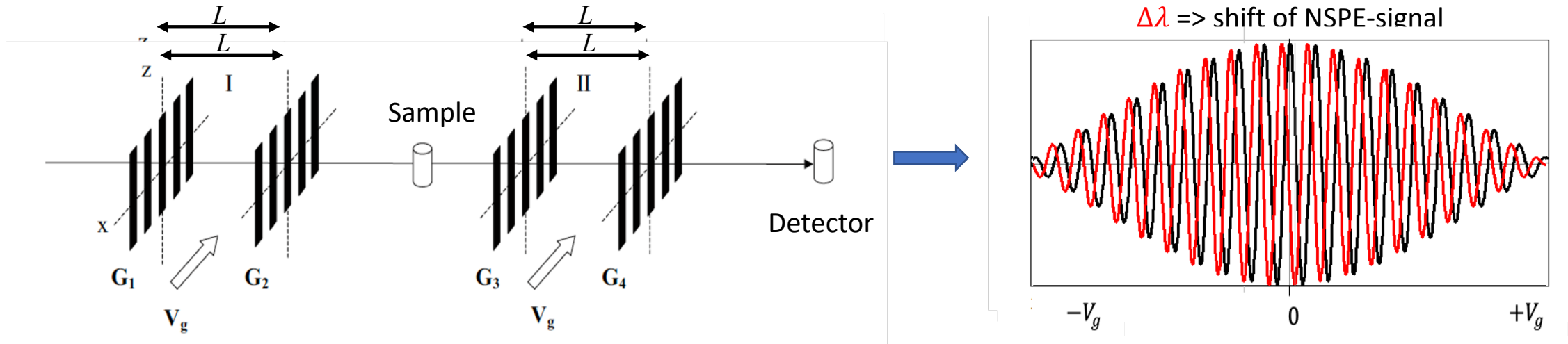


Inelastic scattering on sample:

$\Delta\lambda \Rightarrow$ shift of NSPE-signal



NSPE spectrometer: resolution



We obtain Neutron Spin Echo – like signal, however without the use of neutron spin.

Speed
~~Neutron Spin Echo~~ => no polarization and pol.analysis

Energy resolution:
$$\frac{\Delta E_n}{E_n} = 2 \frac{\Delta v_n}{v_n} = \frac{d}{\pi L} \frac{v_n}{V_g} r$$

r - relative precision of determining the phase shift ($\approx 1-2\%$)

Energy transfer:
$$\Delta E_n = \frac{h^3}{2\pi m_n^2} \frac{d}{L} \frac{1}{V_g} r \lambda^{-3}$$

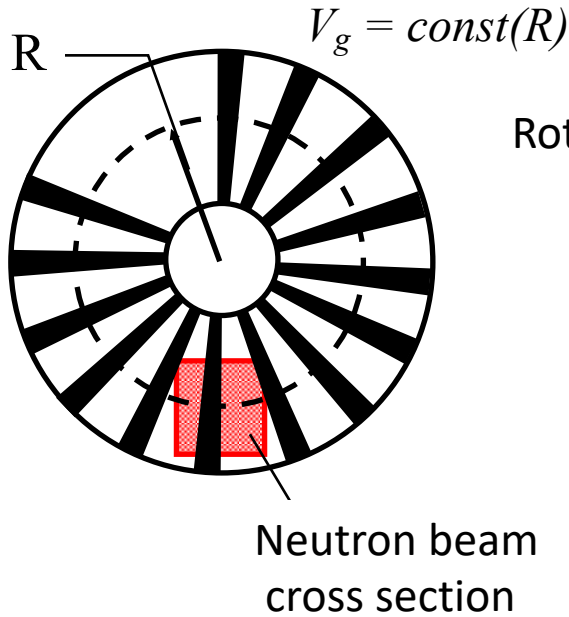
Measurable ΔE_n :

$\lambda=10\text{\AA} \Rightarrow \Delta E_n/E_n \approx 5 \cdot 10^{-7}, \Delta E_n=0.4 \text{ neV}$

Similar to NSE

VCN Speed-Echo spectrometer: search for q_n

Radial diffraction grating



Rotating disk (standard chopper):

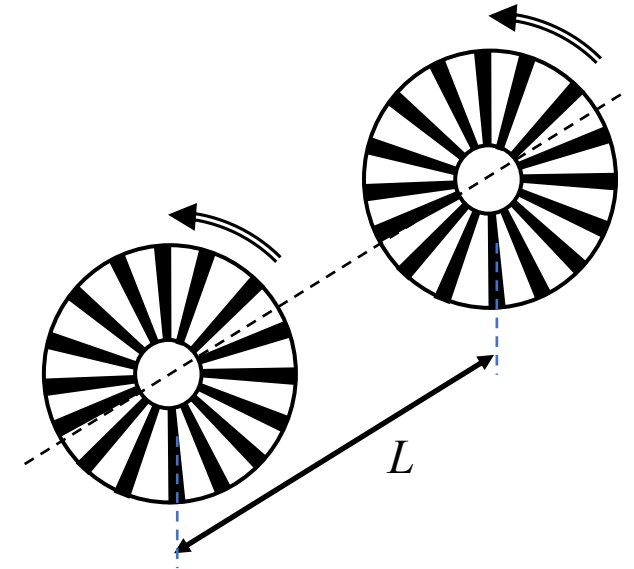
- up to 18,000 rpm
- frequency $f = 300$ Hz

$$f = 300 \text{ Hz}$$

$$R = 30 \text{ cm}$$

$$\downarrow$$

$$V_g = 565 \text{ m/s}$$



$$\Delta E_n = \frac{h^3}{2\pi m_n^2} \frac{d}{L} \frac{1}{V_g} r \lambda^{-3}$$

$$d = 3 \mu\text{m}$$

$$L = 2 \text{ m}$$

$$\lambda = 200 \text{ \AA} \Rightarrow \Delta E_n \geq 3.4 \cdot 10^{-16} \text{ eV}$$

Neutron charge in electric field:

$$\Delta E_n(q_n) = q_n E_{el} L_{el}$$

$$E_{el} = 60 \text{ kV/cm}$$

$$L_{el} = 2 \text{ m}$$

$$\Delta E_n(10^{-20} e) = 1.2 \cdot 10^{-14} \text{ eV}$$

30 times better,
than present day limit

Conclusion

- Interferometry of Very Cold Neutrons (VCN), relies on diffraction gratings for effectively splitting neutron waves coherently.
- The symmetric 4-grating neutron interferometer fully compensates for aberrations from diffraction gratings and Earth's gravity. This type of interferometer functions independent of source coherence, **accommodating non-monochromatic and non-collimated neutron beams.**
- It holds potential for significantly improving the current experimental limit on neutron charge by two orders of magnitude, down to $3 \times 10^{-23}e$. The utilization of holographic (interference) gratings with sub-micrometer periods enhances the performance.
- VCN finds application in Neutron Speed-Echo spectroscopy, facilitating the measurement of extremely small changes in neutron energy. This method also shows potential for investigating neutron charges down to $10^{-23}e$.

Thank you for attention!