

Is it possible to design a high-sensitivity search for neutron-antineutron oscillations at the new UCN/VCN source in Kazakhstan?

- The "clock-non-reset" method,
- Advantages of the new UCN/VCN source in Kazakhstan and the sensitivity estimations,
- Requirements for the infrastructure needed.

- V.V. N., V. Gudkov, K.V. Protasov, W.M. Snow, A.Yu. Voronin, "Experimental approach to search for free neutron-antineutron oscillations based on coherent neutron and anti-neutron mirror reflection", **Phys. Rev. Lett. 122 (2019)** 221801;
- V.V. N., V. Gudkov, K.V. Protasov, W.M. Snow, A.Yu. Voronin, "Comment on B.O. Kerbikov "The effect of collisions with the wall on neutron-antineutron transitions", **Phys. Lett. B 803 (2020)** 135357;
- V.V. N., V. Gudkov, E. Kupryanova, K. Protasov, M. Snow, A. Voronin, "Implementation of neutron/antineutron guides in experiments searching for n - \bar{n} oscillations", **PANIC2021 (2021)** 42B;
- V.V. N., "Why very cold neutrons could be useful for neutron-antineutron oscillation searches", **J. Neutron Res. 24 (2022)** 223.

$$n - \bar{n} \quad \Delta B = 2$$

An **observation** of neutron-antineutron oscillations, which violate both Baryon and Baryon-Lepton conservation, would constitute a scientific discovery of fundamental importance to **physics and cosmology**.

A stringent **upper bound** on its transition rate would make an important contribution to our understanding of the Baryon asymmetry of the universe by eliminating the **post-sphaleron baryogenesis** scenario in the light quark sector.

1. $n - \bar{n}$ oscillations of neutrons in the so-called **quasi-free** limit, when oscillations are not suppressed by external fields (magnetic field, optical potential of residual gases, etc), and thus the probability of oscillations is proportional to the **square of the observation time** (time intervals shorter than $\sim \Delta E / \hbar$);
2. $n - \bar{n}$ oscillations of neutrons **bound in nuclei** (much larger number of neutrons available but much shorter observation times because of the suppression of oscillations by strong nuclei fields).

In any case, the **appearance of antineutrons** is the signature of the process.

At present, both methods provide comparable constraints for the characteristic oscillation time equal to **$\sim 10^8$ sec** (nuclei constraints are more sensitive but model-dependent).

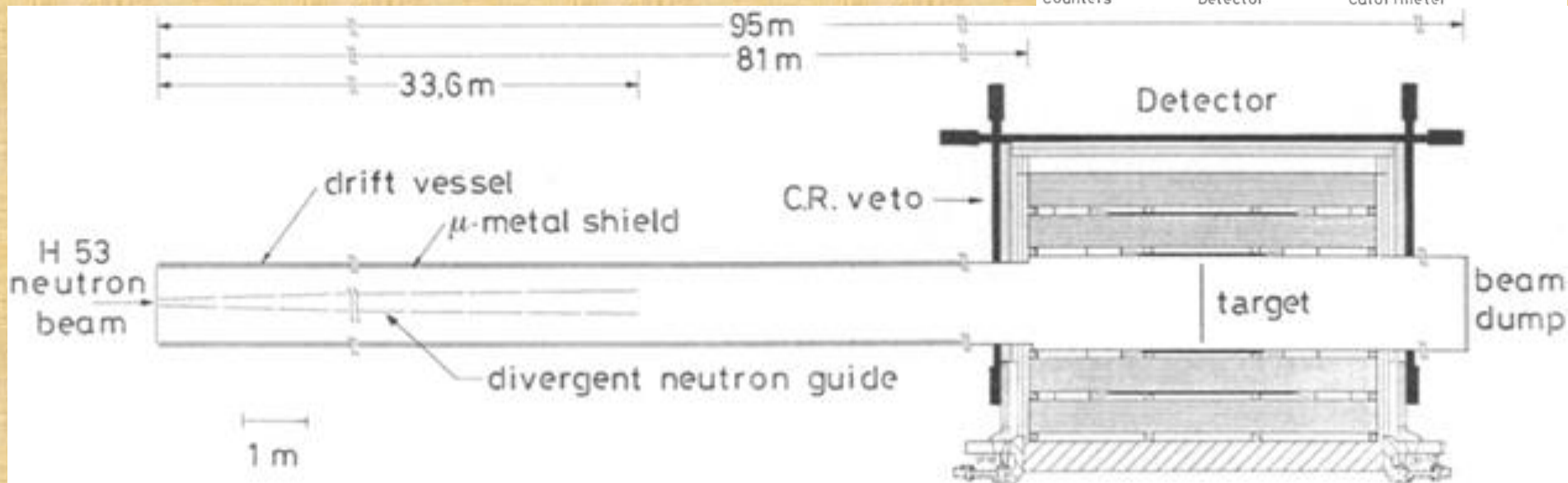
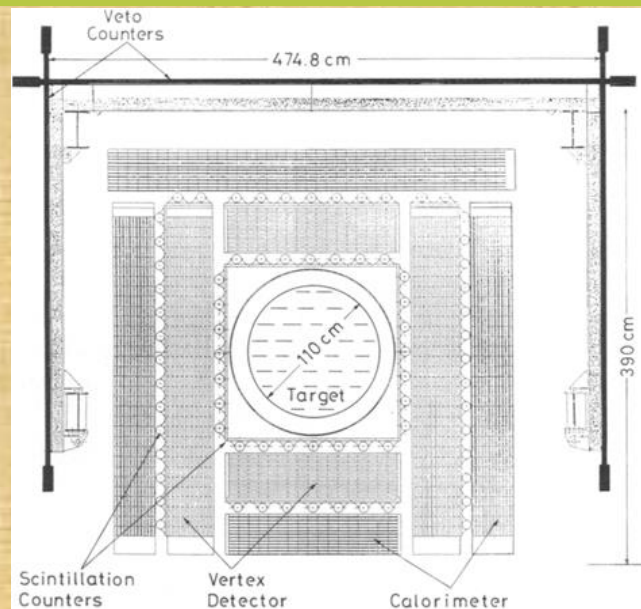
We proposed a **new method**, which combines somehow the advantages of the two methods (**the knowledge of nuclear suppression** of oscillations and (quasi)-**model-free interpretation** of results) and provides a great improvement in the sensitivity (compared to the best existing result).

M. Baldo-Ceolin, et al, "A new experimental limit on neutron-antineutron oscillations", *Z. Phys. C* **63** (1994) 409

$$\tau_{n\bar{n}} > 0.86 \cdot 10^8 \text{ s}$$

Parameters of this experiment:

$$F_n \sim 10^{11} \text{ s}^{-1}; \tau \sim 0.1 \text{ s}; T \sim 2.4 \cdot 10^7 \text{ s}$$



Development of the quasi-free-neutron method: **neutrons** are allowed to **bounce** from the neutron guide walls. An antineutron would travel along the same trajectory, without annihilating and/or losing coherence of the two states for an **extended period of time**.

Analogy to the proposed earlier experiments with **ultracold neutrons** [M.V. Kazarnovski *et al*, JETP Lett. 32 (1980) 82; K.G. Chetyrkin *et al*, Phys. Lett. B 99 (1981) 358; H. Yoshiki, R. Golub, Nucl. Phys. A 501 (1989) 869] but those proposals did not consider coherence of neutrons and antineutrons at reflection, or did not identified conditions at which coherence is maintained.

We:

- Extend this approach to **higher neutron energies**, thus largely increasing statistics and sensitivity,
- Point out conditions for suppressing the **phase difference** for neutrons and antineutrons at reflection,
- Underline the importance of setting **low transverse momenta** of neutrons,
- and making **certain choices for the nuclei** composing the guide material.

For **the same installation length**, include

- **Smaller transversal sizes**,
- **Lower costs**,
- **Larger statistics** (higher accuracy) (can use VCN spectrum).

For **a larger length**,

the **gain in sensitivity**, in terms of the oscillation probability, increases quadratically with length (and still a large reductions of costs)

Our proposal is justified by:

1. (mainly) The validity of quantum mechanics [E. Fermi, *Sul moto dei neutroni nelle sostanze idrogenate*, **Ric. Sci.** **7** (1936) 13.] - the effective optical potential (Fermi potential) is introduced in precisely the same way for neutrons and antineutrons, and
2. (to a smaller extend) some knowledge of antineutron-nuclei scattering lengths [K.V. Protasov, V. Gudkov, E.A. Kupryanova, V.V. N., W.M. Snow, A.Y. Voronin, "Theoretical analysis of antineutron-nucleus data needed for antineutron mirrors in neutron-antineutron oscillation experiments", **Phys. Rev. D** **102** (2020) 075025]

$P_{n \rightarrow \bar{n}} \approx \varepsilon^2 e^{-\frac{\Gamma_a}{2}t} t^2$: The **probability** of neutron-antineutron oscillation depends essentially only on a few parameters: ε , the neutron-antineutron **mixing** parameter, Γ_a , the antineutron annihilation **width**, and time t .

For the optimum observation time $t = \frac{4}{\Gamma_a}$ (obtained by differentiation of the formula above), the probability is:

$P_{n \rightarrow \bar{n}} \approx 2.1 \left(\frac{\varepsilon}{\Gamma_a} \right)^2$ (the optimum is not sharp, and there is only one parameter left).

Crucial **parameters** for the analysis of this problem are:

- The probability of neutron and antineutron reflection per wall collision, ρ_n and $\rho_{\bar{n}}$,
- The difference of phase shifts of the wave function per wall collision, $\Delta\varphi_{n\bar{n}} = \varphi_n - \varphi_{\bar{n}}$.

They depend on:

- The optical potential for neutrons $U_n = V_n + iW_n$, and
- The optical potential for antineutrons $U_{\bar{n}} = V_{\bar{n}} + iW_{\bar{n}}$.

In order to optimize the **sensitivity** of neutron-antineutron searches and simultaneously to decrease the **impact** of theoretical uncertainties, we will use the following limit:

$e \ll V_n, e \ll V_{\bar{n}}, e \sim W_{\bar{n}}, W_n \ll V_n, W_{\bar{n}} \ll V_{\bar{n}}, W_n \ll W_{\bar{n}}$, with e the energy of transversal neutron motion. Then,

for the probabilities: $\rho_n = 1$ and $1 - \rho_{\bar{n}} \approx \frac{2kk_{\bar{n}}''}{(k'_{\bar{n}})^2}$, with

$k'_{\bar{n}} \approx \sqrt{2mV_{\bar{n}}}$ and $k_{\bar{n}}'' \approx \sqrt{m \left(\frac{W_{\bar{n}}^2}{2V_{\bar{n}}} \right)}$ and for the phase

shift: $\Delta\varphi_{n\bar{n}} \approx \frac{2k}{k_n k'_{\bar{n}}} (k_n - k'_{\bar{n}})$

Imagine two upstream sections a two-dimensional ballistic neutron guide (with a cross-section increasing from h by d to H by D). Typical cross-sections are $hd \sim 10^2 \text{ cm}^2$, $HD \sim 10^4 \text{ cm}^2$, respectively. In accordance with Liouville theorem, tangential velocity components would decrease from $\sim 2v_{crit}^{Ni}$ to

$$|v_{hor}| < 2v_{crit}^{Ni} \frac{d}{D} \text{ and } |v_{vert}| < \sqrt[3]{4hdv_{crit}^{Ni}g}.$$

$$b_{\bar{n}A} \sim 1.54 \sqrt[3]{A} - i$$

Element	$b_{\bar{n}A}$ [fm]	$U_{\bar{n}}$ [neV]	$\tau_{\bar{n}}$ [s]
C	3.5 - i	103 - $i29$	1.7
Mg	3.5 - i	39 - $i11$	1.0
Si	3.7 - i	48 - $i13$	1.2
Ni	4.7 - i	111 - $i24$	2.3
Cu	4.7 - i	104 - $i22$	2.2
Zr	5.3 - i	59 - $i11$	1.8
Mo	5.3 - i	89 - $i16$	2.3
W	6.5 - i	106 - $i16$	3.0
Pb	6.7 - i	57 - $i8.6$	2.3
Bi	6.7 - i	49 - $i7$	2.1

Then, $\tau_{hor}^{\Delta\varphi, \bar{n}} = \frac{D}{|v_{hor}|} \cdot \frac{\sqrt{V_n V_{\bar{n}}}}{2\sqrt{e_{hor}}(\sqrt{V_n} - \sqrt{V_{\bar{n}}})} \sim 32 \text{ s}$ and

$$\tau_{vert}^{\Delta\varphi, \bar{n}} = \frac{|v_{vert}|}{g} \frac{\sqrt{V_n V_{\bar{n}}}}{\sqrt{e_{vert}}(\sqrt{V_n} - \sqrt{V_{\bar{n}}})} \sim 7.3 \text{ s}$$

Note, however, that a factor $\left((\sqrt{V_n} - \sqrt{V_{\bar{n}}}) \rightarrow 0 \right)$ can allow to largely increase these characteristic times by proper mixing of two isotopes/elements for the guide wall material if needed.

$$\tau_{hor}^{\rho, \bar{n}} = \frac{D}{|v_{hor}|} \frac{(V_n)^{3/2}}{W_{\bar{n}} \sqrt{e_{hor}}} \sim 15 \text{ s},$$

$$\tau_{vert}^{\rho, \bar{n}} = \frac{2|v_{vert}|}{g} \frac{(V_n)^{3/2}}{W_{\bar{n}} \sqrt{e_{vert}}} \sim 3.1 \text{ s}.$$

$\tau_{vert}^{\rho, \bar{n}}$ is THE real limitation of this method.

Even in the limit of "zero" vertical velocities, this estimation will not significantly change.

You can improve this value by using a vertical extraction of neutrons, then nearly no constraints on the observation time !

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INP Kazakhstan:

VCN phase-space density 3-5 \downarrow , VCN velocity 20 \downarrow , thus,
 $F_n \sim 10^9 s^{-1}$

$$\tau \sim 3.0 s ; T \sim 2.2 \cdot 10^8 s$$

Thus, the gain factor is **~80**

Additional gain factors to be discussed:

- With an advanced VCN extraction system,
~10-20
- With the vertical neutron guide of **~100 m**,
an additional gain factor of **~10**

- A **VCN source** (for instance, a helium VCN source surrounded by a layer of fluorinated nanodiamonds, probably it might be **combined with the UCN source**),
- An extraction system which provides the **largest extracted VCN flux** (**simplified by the geometry of the thermal column**),
- A **horizontal VCN guide** with the length of **200-400 m** (**is there sufficient space available around the reactor?**),
- Long measuring time **~10 years** (**would the reactor still be running?**),
- A significant **budget** (**an international collaboration?**).